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A LEVEL-SET METHOD FOR MAGNETIC SUBSTANCE SIMULATION

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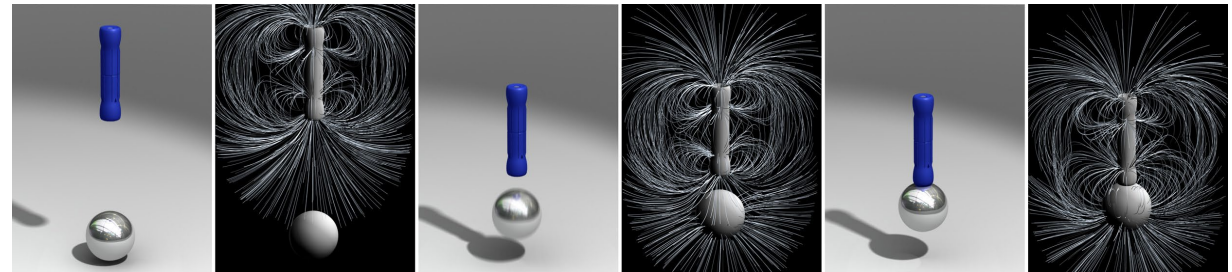
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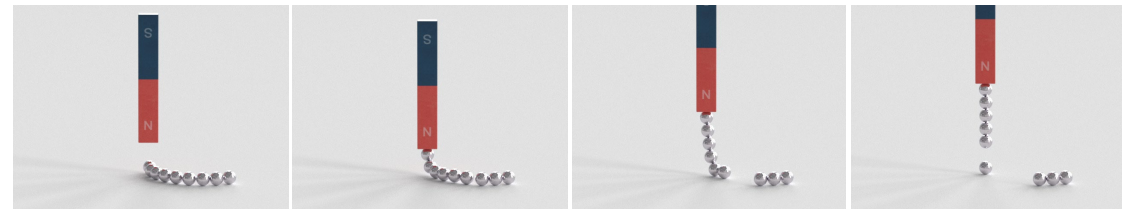
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MAGNETIC SIMULATION

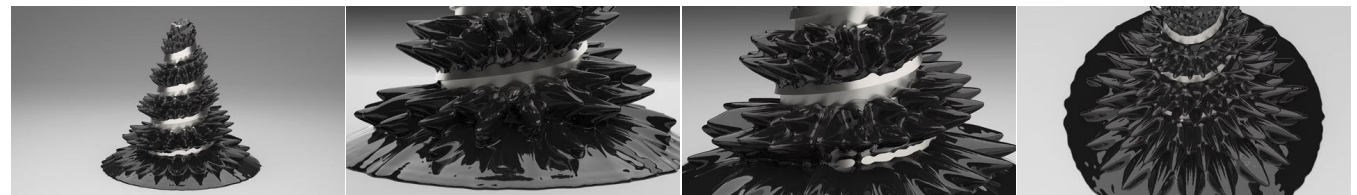
- Magnetic interaction
 - Phenomena in life
 - Applications in industry
- Magnetic simulation in graphics
 - Magnetic rigid bodies
 - [Thomaszewski et al. 2008]
 - [Kim et al. 2018]
 - Magnetic liquids (ferrofluids)
 - [Huang et al. 2019]
- Observation
 - Need for a unified framework
 - All based on Lagrangian views



[SIGGRAPH 2008] Magnets in Motion



[SIGGRAPH 2018] Magnetization Dynamics for Magnetic Object Interactions

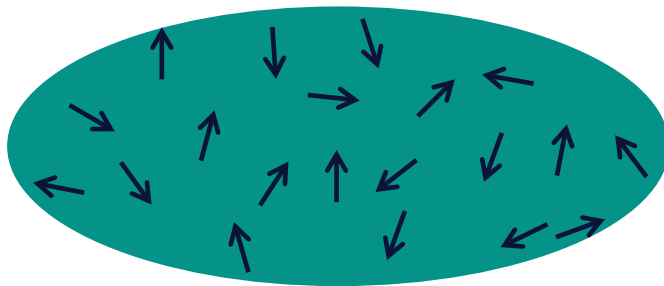


[SIGGRAPH 2019] On the Accurate Large-scale Simulation of Ferrofluids

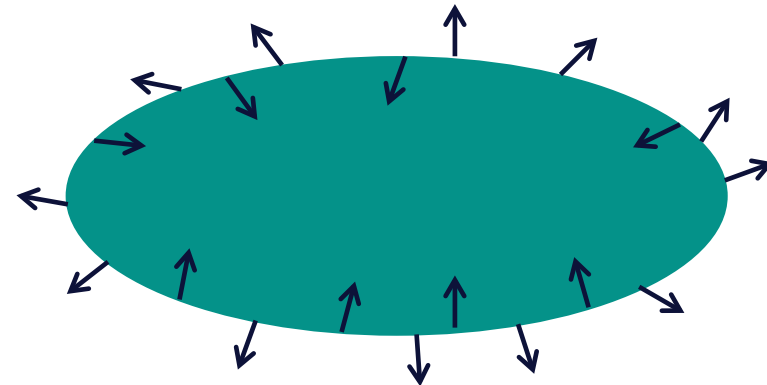
- Treat magnetic interaction as interfacial force
 - Solving the magnetic-mechanical coupling problem efficiently
- Based on an Eulerian view
 - Taking advantage of the level-set method
- Model a broad range of magnetic phenomena in a unified way
 - Treating classical simulators simply as black boxes
 - Supporting fluids, solids, and their couplings
- Include an efficient, precise numerical scheme
 - Introducing magnetic force by solving a Poisson equation with jump conditions
 - Especially easy to be incorporated into a standard Euler fluid solver

HOW TO PORTRAY MAGNETIC FORCE

- The first step to simulate the motion of a physical system
 - To establish the dynamics of this system
 - Equivalent to determining what the form of force is in Newtonian mechanics
- Magnetic force:
 - Yet to be definitely answered



Volumetric (throughout the substance)



Interfacial (only on the surface)

HOW TO PORTRAY MAGNETIC FORCE

- The Maxwell stress tensor \mathbf{T}_m
 - $\mathbf{T}_m = \frac{1}{\mu_0} \left(\mathbf{B} \otimes \mathbf{B} - \frac{1}{2} B^2 \mathbf{I} \right)$ in vacuum, with the electric terms omitted
 - μ_0 is the constant *vacuum permeability*
 - \mathbf{B} is the *magnetic induction intensity*
 - \mathbf{I} is the second-order unit tensor
 - Physicists compute magnetic force by taking $\mathbf{f}_m = \nabla \cdot \mathbf{T}_m$
- Abraham–Minkowski controversy
 - What the form of \mathbf{T}_m is in matter
 - **Einstein–Laub form** vs. **Minkowski form**

Form	The Maxwell stress tensor	Magnetic force
Einstein–Laub	$\mathbf{T}_m^E = \mathbf{B} \otimes \mathbf{H} - \frac{\mu_0}{2} H^2 \mathbf{I}$	$\mathbf{f}_m^E = \mu_0 \mathbf{M} \cdot \nabla \mathbf{H}$
Minkowski	$\mathbf{T}_m^M = \mathbf{B} \otimes \mathbf{H} - \frac{1}{2} (\mathbf{B} \cdot \mathbf{H}) \mathbf{I}$	$\mathbf{f}_m^M = \mathbf{B} \cdot \nabla \mathbf{H} - \frac{1}{2} \nabla (\mathbf{B} \cdot \mathbf{H})$

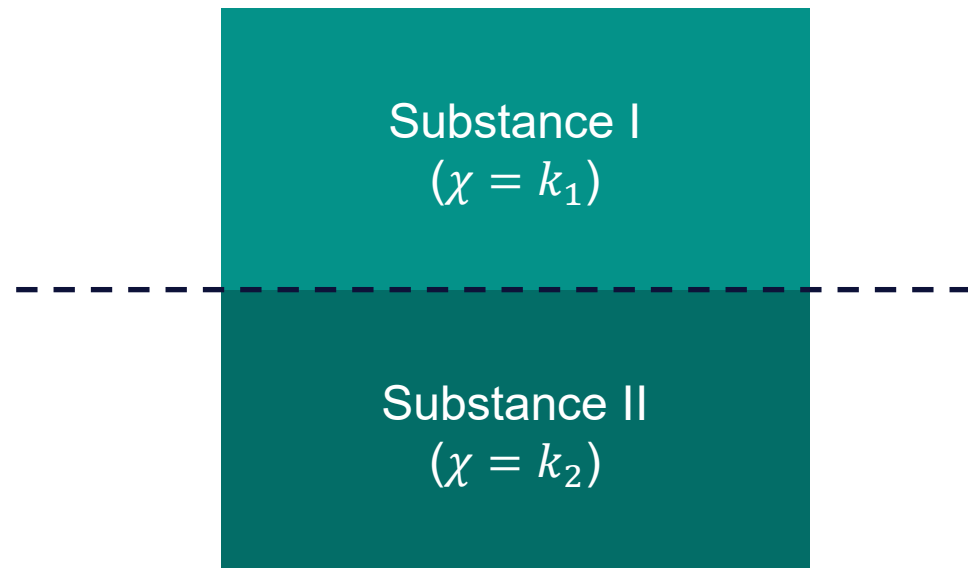
* \mathbf{H} is the *magnetic field intensity* and \mathbf{M} is the *magnetization intensity*, satisfying $\mathbf{B} = \mu_0(\mathbf{H} + \mathbf{M})$

\mathbf{f}_m^E is called the *Kelvin force*,
while \mathbf{f}_m^M is called the *Helmholtz force*

Totally Different!

HOW TO PORTRAY MAGNETIC FORCE

- Assume that the magnetic substances in the system are all linear, isotropic: $\mathbf{M} = \chi \mathbf{H}$
 - χ is called the *magnetic susceptibility*, constant within each substance
 - Kelvin force vs. Helmholtz force
 - $\mathbf{f}_m^E = \frac{\mu_0}{2} \chi \nabla(H^2)$
 - Compatible with the Lagrangian view
 - Let $\chi_i = k_1$ if Particle i belongs to Substance I and $\chi_i = k_2$ if Particle i belongs to Substance II
 - The magnetic force exerted on Particle i is $\frac{\mu_0}{2} \chi_i \nabla(H^2)$
 - These are what related works adopted
 - $\mathbf{f}_m^M = -\frac{\mu_0}{2} H^2 \nabla \chi$
 - Only on the interface!



* k_1 and k_2 are both constants

HOW TO PORTRAY MAGNETIC FORCE

Question: since the two forms are totally different, can they both be right?

- Take fluid simulation as an example
 - Need solving the pressure to enforce incompressibility
 - $\nabla \cdot [\rho u + (\mathbf{f} - \nabla p)dt] = 0$
 - ρu is the advected momentum density
 - BC: $p = 0$ outside the fluid
 - Provided that there is an external force \mathbf{f}'
 - $\mathbf{f}' = \nabla \Phi$, satisfying $\Phi = 0$ outside the fluid
 - Theorem: \mathbf{f}' does not influence the motion
 - $p^* = p + \Phi$ is a new solution of the pressure
 - $\nabla \cdot [\rho u + (\mathbf{f} + \mathbf{f}' - \nabla p^*)dt] = 0$
 - $p^* = 0$, outside the fluid
 - The total force leaves the same
 - $\mathbf{f} - \nabla p = \mathbf{f} + \mathbf{f}' - \nabla p^*$

$$\mathbf{f}' = \mathbf{f}_m^E - \mathbf{f}_m^M = \nabla \left(\frac{\chi \mu_0 H^2}{2} \right)$$

- \mathbf{f}_m^E and \mathbf{f}_m^M are indistinguishable in effect
- The essence of the A–M controversy
 - How to divide physical quantities between ferromagnetic system and mechanical system

Ferromagnetic
force

Undecided
 \mathbf{f}'

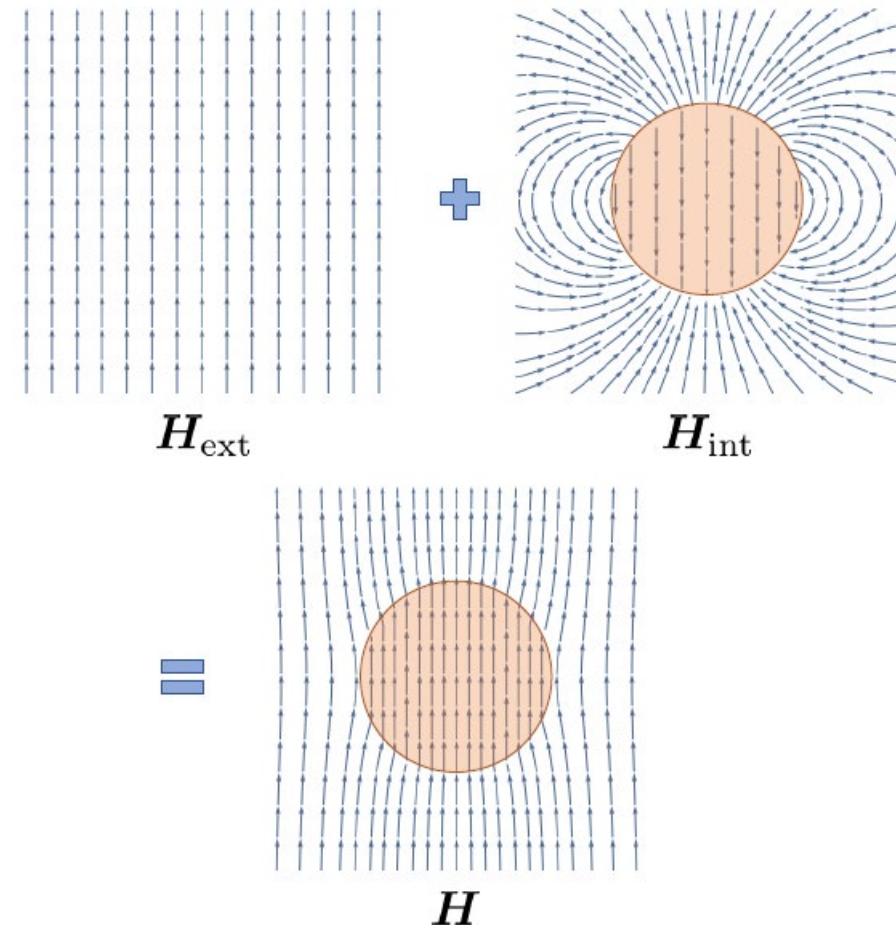
Mechanical
force

- Eulerian view & Helmholtz force
 - Capture the surface accurately by the level-set
 - Reduce \mathbf{f}_m from 3D to 2D
- Four-step interaction
 - Magnetization
 - \mathbf{H}_{ext} magnetizes the immersed magnetic substance
 - According to its current shape and position
 - Induction
 - \mathbf{H}_{int} is induced
 - $\mathbf{H} = \mathbf{H}_{\text{ext}} + \mathbf{H}_{\text{int}}$
 - Exertion
 - Apply the magnetic force $\mathbf{f}_m = \mathbf{f}_m^M$
 - Reshaping
 - The state of the physical system is updated

* Simultaneously and repeatedly

Time Splitting

- Magnetization & induction
 - Solve for \mathbf{H}_{int} based on \mathbf{H}_{ext} and χ
 - Governed by Maxwell's equations
 - $\nabla \cdot \mathbf{B} = 0$
 - $\nabla \times \mathbf{H} = \mathbf{j}_f + \frac{\partial \mathbf{D}}{\partial t}$
 - Equivalent to a Poisson's equation
 - $\nabla \cdot (1 + \chi) \nabla \psi = \nabla \cdot \chi \mathbf{H}_{\text{ext}}$
 - $\mathbf{H}_{\text{int}} = -\nabla \psi$
 - (Discretized into a linear system)



- Exertion

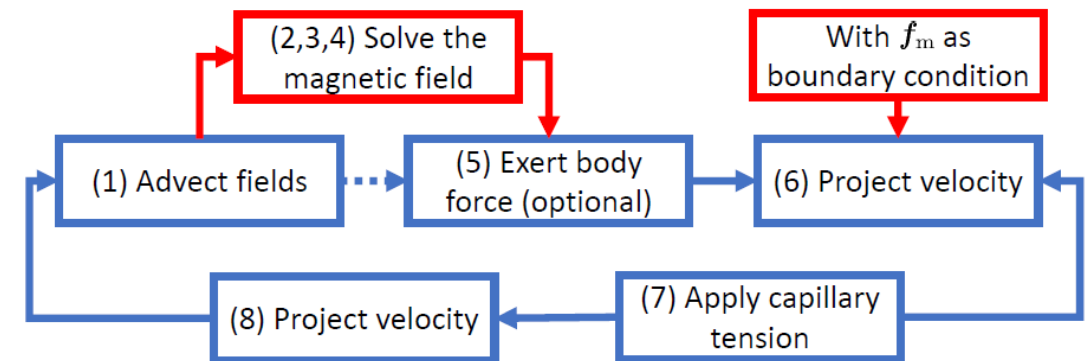
- The precise formula of **Helmholtz force**
 - $\mathbf{f}_m = \frac{\mu_0}{2} k \left[\mathbf{H}^2 + \frac{k^2}{4k+4} (\mathbf{H} \cdot \hat{\mathbf{n}})^2 \right] \delta_\Sigma(\mathbf{r}) \hat{\mathbf{n}}$
 - δ_Σ is the generalized *Dirac delta function*
 - $\hat{\mathbf{n}}$ is the unit normal vector on the interface
 - Derived by taking the weak form of δ_Σ
 - Refer to the paper for details
- Always perpendicular to the interface

- Reshaping

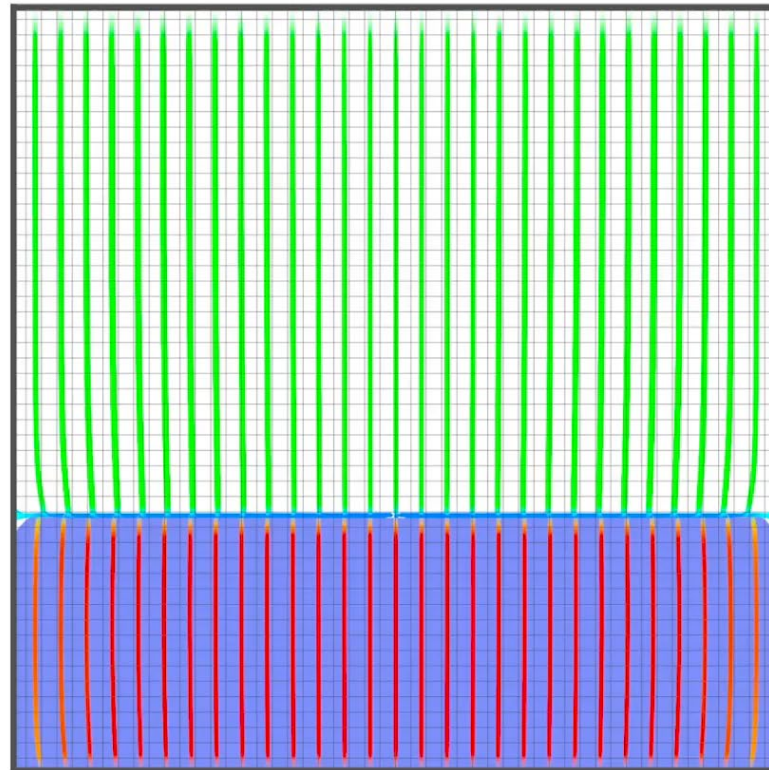
- For fluids (with surface tension)
 - $\rho \left(\frac{\partial \mathbf{u}}{\partial t} + \mathbf{u} \cdot \nabla \mathbf{u} \right) = -\nabla p + \rho \mathbf{g} + \mathbf{f}_m - \sigma \kappa \hat{\mathbf{n}}$
 - $\nabla \cdot \mathbf{u} = 0$
- For elastic bodies
 - $\rho \left(\frac{\partial \mathbf{u}}{\partial t} + \mathbf{u} \cdot \nabla \mathbf{u} \right) = \nabla \cdot \boldsymbol{\sigma} + \rho \mathbf{g} + \mathbf{f}_m$
- For rigid bodies
 - $m_\rho \frac{d\mathbf{u}}{dt} = m_\rho \mathbf{g} + \iiint_\Omega \mathbf{f}_m dV$
 - $\mathbf{I}_\rho \frac{D\boldsymbol{\omega}}{Dt} + \boldsymbol{\omega} \times (\mathbf{I}_\rho \cdot \boldsymbol{\omega}) = \iiint_\Omega (\mathbf{d} \times \mathbf{f}_m) dV$

- Solve the magnetic field (a linear system)
 - Conjugate gradient method
 - Multi-grid preconditioned
- Exert the magnetic force
 - For fluids
 - Treat the interfacial **Helmholtz force** as pressure jump
 - Young–Laplace Equation
 - $f_m dV = \Delta p ds$
 - Involved as BC in projection step
 - For solids
 - Apply on surface triangles

* The pipeline sketch of ferrofluid simulation with new components painted red



SIMULATION RESULTS



Rosensweig instability (2D)

CONCLUSION

- Summary
 - The first versatile level-set method
 - A novel computational approach
 - With interfacial **Helmholtz force** introduced
 - An efficient numerical scheme
- Limitations
 - Explicit magnetic force
 - Potential numerical instability
 - Weak coupling
- Future work
 - Boundary element method
 - Apply Helmholtz force implicitly

Thank you for listening!